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CEELS as a Probe of the Carbide to Graphite Transformation on NI

Ву

DAVID E. RAMAKER



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# CEELS AS A PROBE OF THE CARBIDE TO GRAPHITE TRANSFORMATION ON Ni

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#### **ABSTRACT**

Recent effective medium theory calculations indicate that a nucleation step must be involved in the carbide to graphite transformation on Ni, and that vertical  $C_2$  may be a precursor to this nucleation step. A similar nucleation site could also be involved in diamond formation on metals. We interpret previously published C K edge CEELS and NEXAFS data for carbon/Ni between 500-700K. We make comparison with theoretical DOS calculations, and utilize CEELS angle dependencies to assign features and determine bond orientations on the surface. The results confirm that horizontal  $C_n$ , most likely  $C_2$ , exists on the surface up 600K, but that at 620K, some of the  $C_2$  species flip up. Evidence for vertical  $C_2$  is seen only at higher C coverages suggesting that indeed these vertical  $C_2$ 's serve as the precursor for the nucleation of graphite, and perhaps also for diamond.

#### THE CARBIDE TO GRAPHITE TRANSFORMATION

Carbon is relatively unreactive with Ni. Thus Ni forms a relatively unstable carbide, which above 670K either undergoes dissolution into the bulk or at higher coverages forms a graphitic layer which sits high above the surface (e.g. 2.8 Å above a Ni(111) surface) [1]. Recent calculations utilizing effective medium theory (EMT) [1] shed some light on this transformation. In the EMT approach, the atom positions are determined by the electron density, each atom seeking its own unique optimum density. These calculations conclude that at intermediate coverage, the C-C interaction drives carbon closer to a Ni surface, but in a graphite layer, the C-C interaction drives carbon away from the surface. They conclude that the carbide to graphite evolution is discontinuous, indicating that a nucleation step must be involved.

The calculations also suggest that horizontal  $C_2$  species on a Ni(111) surface are not stable at high temperature, but that vertical (i.e. perpendicular)  $C_2$  species may be stable on the Ni surface [1]. However, this vertical  $C_2$  is too close to the surface to serve as a nucleation site for graphite formation by itself. A  $C_3$  species is the smallest entity which moves sufficiently far from the Ni surface to serve as a graphite nucleation site. Darling et al. [1] then postulate that a vertical  $C_2$  species may either "tip over" to form a  $C_3$  species, which moves

away from the surface to act as a graphite nucleation site, or a graphitic layer forms on top of a carbidic layer (i.e. the outer C of the vertical  $C_2$  may become part of the graphite layer, and the inner C may ultimately undergo dissolution into the bulk). In either case, a vertical  $C_2$  serves as a precursor to the nucleation site.

An identical nucleation step could be involved in the nucleation of diamond on metal surfaces, and critically determine the bonding of the diamond film to the metal. Thus we search for experimental confirmation of these theoretical results for graphite. NEXAFS (near edge x-ray absorption fine structure) data have been very helpful in determining the various orientations of molecular hydrocarbon fragments on Ni below 450 K [2]. but little new information from NEXAFS data has been obtained for C/Ni above 450 K. In this work we utilize previously published [3,4] CEELS (core electron energy loss spectroscopy) and NEXAFS data [5] to obtain experimental verification that vertical C<sub>2</sub> acts as a precursor for graphite nucleation on Ni.

Recently Caputi et al. [3] reported AES and CEELS data for carbon on Ni(100) in the range 520-770 K. Core-level and valence band XPS data have also been reported [3]. A detailed interpretation of the AES and XPS data has revealed extensive C-C bonding on the surface, in what was previously believed to be carbidic (i.e. only C-M bonding) in character [6]. At higher temperatures, just prior to the formation of graphite around 620K, the AES data indicate that the amount of C-C bonding appears to decrease. No meaningful interpretation of the CEELS data in this same temperature range has been reported.

#### CEELS ANGULAR DEPENDENCE

In NEXAFS, the dipole selection rule (i.e.  $s \rightarrow p$  only) is appropriate. By CEELS, we mean the use of small electron energies (500-1000 eV) and the measurement of back scattered electrons which have suffered large momentum transfer, in which case the validity of the dipole selection rule is not expected. Thus optically forbidden monopole transitions should be evident. Nevertheless, CEELS data can still be utilized to obtain some of the same information obtainable from NEXAFS data.

We utilize equations derived by Cheung [7] for determining the angular dependence of CEELS data for graphite, which has the  $\sigma$  orbital parallel to the surface. Cheung obtains,

$$\{1s \to \pi\} \propto 1.5 \xi \sin^2 \delta + 3 \nu [\cos^2 \delta - 0.5 \sin^2 \delta] \sin^2 \theta \qquad (1)$$

$$\{1s \to \sigma\} \propto \varepsilon/3 + \nu[1-0.5\sin^2 \varepsilon] - \nu[\cos^2 \varepsilon - 0.5\sin^2 \varepsilon] \sin^2 \theta, \qquad (2)$$

where  $\delta$  is the electron acceptance angle and  $\xi$  and  $\nu$  are the magnitude of the monopole and dipole contributions, respectively. Here  $\theta$  is the angle between

the surface and the electron beam direction. We compare these expressions with those appropriate for NEXAFS and C=C bonds parallel to the surface [2],

$$\{1s \to \pi\} \propto 6\cos^2\theta \tag{3}$$

$$\{1s \to \sigma\} \propto 3\sin^2\theta \tag{4}$$

$$\{1s \to \sigma\} \quad \propto \quad 3\sin^2\theta \,. \tag{4}$$

If we assume that  $\delta = 90^{\circ}$  in eqs. 1 and 2 (i.e. that electrons are counted at all acceptance angles), the CEELS technique mimics the NEXAFS technique, which utilizes the total electron yield or the Auger yield. We then obtain,

$$\{1s \to \pi\} \propto 1.5\nu \cos^2\theta \qquad (5)$$

$$\{1s \to \sigma\} \propto \xi/3 + 0.5\nu + 0.5\nu \sin^2\theta. \qquad (6)$$

$$\{1s \to \sigma\} \propto \xi/3 + 0.5\nu + 0.5\nu \sin^2\theta. \tag{6}$$

These expressions have similar  $\sin^2\theta$  and  $\cos^2\theta$  dependences to those above for NEXAFS as expected. In contrast, if we assume  $\delta = 16.5^{\circ}$  (i.e. the appropriate acceptance angle for a cylindrical mirror analyzer and a 500 eV excitation beam as utilized by Cheung [7]), we obtain the expressions,

$$\{1s \rightarrow \pi\} \propto 0.12\nu + 2.64\nu \sin^2\theta$$

$$\{1s \rightarrow \sigma\} \propto \xi/3 + 0.08\nu + 0.88\nu \cos^2\theta.$$
(8)

$$\{1s \to \sigma\} \quad \propto \quad \xi/3 + \quad 0.08\nu \quad + \quad 0.88\nu \cos^2\theta \,. \tag{8}$$

Notice the switch in  $\cos^2\theta / \sin^2\theta$  dependence between eqs. 5, 6 and 7, 8. Eqs. 7 and 8, appropriate for this case, are also different from that for NEXAFS, eqs. 3 and 4. Furthermore, Cheung [7] found empirically for graphite with a 500 eV excitation beam that  $\varepsilon/\nu$  is about 8. Thus in CEELS for  $\theta = 90^{\circ}$  (i.e. electron beam perpendicular to surface, which is generally the case for the data discussed in this work) the  $\pi$  and  $\sigma$  contributions have nearly equal intensity (eqs. 7 and 8 above both give 2.75 $\nu$ ). If the C-C bond is vertical to the surface, we expect a corresponding reversal in the dipole intensity dependencies giving  $\{1s \rightarrow \pi\}$  $0.08\nu \text{ and } \{1s \to \sigma\} = 5.42\nu.$ 

In summary, the CEELS data should give about equal  $\sigma$  and  $\pi$  area intensities for parallel C-C orientation, and be dominated by  $\sigma$  intensity for perpendicular C-C orientation.

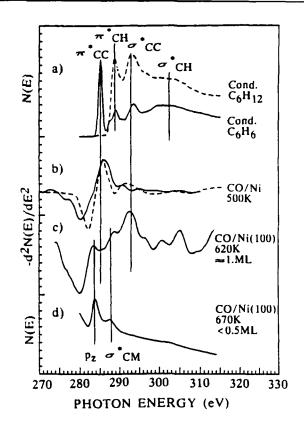
#### INTERPRETATION OF CEELS FEATURES

Figs. 1b and c compare  $-d^2N(E)/dE^2$  K edge CEELS curves for various carbonaceous layers on metals [3,4] along with NEXAFS data for CO/Ni(100) at 670K in 1d [5]. We have also included in Fig. 1a NEXAFS [8] data for condensed benzene and cyclohexane for comparison. NEXAFS data for a large number of molecules and solids show C-H and C-C  $\pi^*$  and  $\sigma^*$  features at surprisingly specific energies [8]; so we have labelled these features accordingly in Fig. 1a.

All of the carbonaceous layers were prepared by exposure of the Ni surfaces to CO at around 500K, and then heating. The dashed curve in 1b was

Fig. 1

- a) C K NEXAFS data for condensed benzene and cyclohexane [5],
- b)  $-d^2N(E)/dE^2$  C K edge CEELS data for CO/Ni(100) at 500K (dotted line) [3] and CO/Ni(111) at 520K (solid line [4]),
- c) CEELS data for CO/Ni(100) at 620K [3], and
- d) NEXAFS data for CO/Ni(100) at 670K, but at low coverage [5].



reported by Rosei et al [4] for a Ni(111) surface upon heating to 500K with an estimated coverage of about 0.3 ML (this is a very crude estimate). The solid curves in b (at 520K) and c (at 620K) were reported by Caputi et al [3]. Although they do not estimate the C coverage, it is believed to be greater than or equal to 1 ML. The NEXAFS curve in (d) was reported by Stohr and Jaeger [5] with incident angle  $\mathfrak{D}=20^{\circ}$ . Consistent with the expected NEXAFS angular dependence, this curve emphasizes vertical  $\mathfrak{D}$  bonds. The latter curve corresponds to less than 0.5 ML of "carbidic" C on the surface. The K binding energy is about 282.9 eV for a carbidic layer on Ni(100) [9]. We have deconvoluted the Caputi data by a 2 eV Gaussian line shape to regain better resolution since it was taken with a large 6  $V_{ptp}$  modulation voltage.

The deconvolved Caputi data on Ni(100) and the Rosei data on Ni(111) at 500 K are quite similar as expected. They reveal the characteristic  $\pi^*_{CC}$  and  $\sigma^*_{CC}$  peaks at 285 eV and 293 eV respectively as seen in the molecular NEXAFS data. Since C-M bonds do not produce peaks in this energy range [10], this clearly indicates the presence of C-C bonding on the surface, consistent with the AES and XPS data [6] as indicated above. The similar area intensities of the  $\pi^*_{CC}$  and  $\sigma^*_{CC}$  peaks indicate that the C-C bonds lie flat on the surface. We would assume that these C-C bonds primarily exist as  $C_n$  (n = 2,3 etc, with n = 2 favored) species on the surface.

The Caputi data at 620K (curve c) reveal dramatic differences from that at 520K. Now the  $\sigma^*_{CC}$  feature dominates with the  $\pi^*_{CC}$  feature nearly missing. This is consistent with C-C bonds perpendicular to the surface. We believe the  $\sigma^*_{CC}$  feature now arises from  $C_2$  species standing erect on the surface.

NEXAFS data for CO/Ni(100) at 300 K [5] (not shown) are dominated by the  $\pi^*$  and  $\sigma^*$  CO bond features. Heating to 670K breaks all C-O bonds, leaving only atomic C on the surface. Notice that in Fig. 2d, no evidence exist for either C-O or C-C bonds. The features at 284 and 288 eV are attributed [10] to nonbonding  $p_z$  orbitals and  $\sigma^*_{CM}$  orbitals bonding the atomic C to the surface. Evidence for these same features also exists in the deconvolved 620K data of Caputi (curve 1c).

#### COMPARISON WITH THEORETICAL DOS

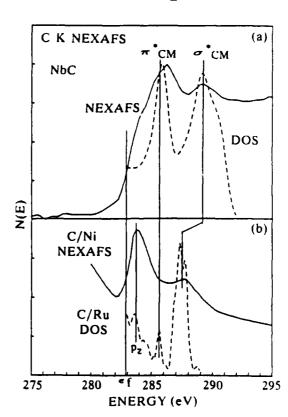
Fig. 2 provides further evidence for our assignment of the p<sub>Z</sub> and  $\sigma^*_{CM}$  features above. We compare NEXAFS data for NbC [11] and C/Ni (i.e. the data in Fig. 1) to DOS results for NbC reported by Schwarz [12] and for C/Ru reported by Feibelman [13]. Schwarz utilized a self-consistent augmented plane wave (APW) X<sub>\infty</sub> calculation for a periodic NbC solid. Feibelman's results were obtained for a (1x1) overlayer of C atoms on an Il-layer Ru(0001) film, with only 1/3 of the three-fold sites filled with carbon atoms. In this theoretical model, the carbon atoms are relatively isolated since the nearest C-C distance is greater than 5.0 a.u. compared with a C-C distance of 2.68 a.u. in graphite [13]. We believe Feibelman's results should be adequate for low coverages of C/Ni since both Ru and Ni form relatively unstable carbides compared with Nb or Ti.

The DOS from both calculations are available for only the first 10 eV or less of the conduction band. Schwarz identified the two peaks in the DOS as the  $\pi$ \* and  $\sigma$ \* antibonding features (i.e. involving the  $t_{2g}$  and  $e_{g}$  Nb d orbitals). Feibelman identified the two major peaks for C/Ni as the  $p_{z}$  and  $\sigma$ \* features.

Fig. 2

Comparison of C K NEXAFS data for NbC and C/Ni with DOS calculations.

- a) NEXAFS data from Wesner et al [11], shifted up in energy by 3.0 eV based on calibration with similar data for graphite. DOS calculation from Schwarz [12].
- b) NEXAFS data from Stohr and Jaeger [5] for CO/Ni at 670K with DOS from Feibelman for a (1x1) layer of C on a Ru(0001) film.



The p<sub>2</sub> feature arises from the C nonbonding or "dangling" bonds pointing outward from the surface. The  $\sigma^*$  feature arises from the three sp<sup>3</sup>-like C-metal bonds per C atom. Note that a similar  $\pi^*$  feature also appears in Feibelman's DOS and aligns with the same peak for NbC, but it has much weaker intensity. Furthermore, note that the  $\sigma^*$  feature has a much lower energy for C/Ni than for NbC. This reflects the much weaker C-metal interaction in Ni-C bonds compared with Nb-C. The comparison with the DOS also helps to identify the precise position of the Fermi level. We place the Fermi level at 282.9 eV.

#### CONCLUSIONS

In summary, our interpretation of the spectroscopic results are consistent with our previous interpretations of the AES and XPS data [6], and with Darling's EMT theoretical results as discussed above [1]. First, the CEELS data do indeed verify that significant horizontal C-C bonding exists on the surface below 600K. We anticipate that this is in the form of C<sub>n</sub>. Around 620K, some vertical  $C_2$  is formed along with considerable  $C_1$  (i.e. atomic C). However, the CEELS and NEXAFS data also suggest that vertical C2 is formed only at higher coverages (it is clearly present in the Caputi data [coverage about 1 ML] but absent in Stohr and Jaeger's NEXAFS data [coverage < 0.5 ML]). At higher coverages, we envisage that some of the C2 are forced to flip up to make room for the neighboring horizontal C2's to dissociate. These vertical C2's may then serve as the precursor for the nucleation of graphite, since it would explain the lack of graphite formation from  $C_2H_4/Ni(100)$  (i.e. at low carbon coverages). Further exposure to  $C_2H_2$  at higher temperatures (i.e. producing higher C coverages) does lead to graphite formation on Ni(100) [9]. Thus we provide the first experimental evidence for vertical C<sub>2</sub> as a precursor for graphite nucleation and corroborate the theoretical EMT [1] results.

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